Modeling of Plasma Rotation and Density Modification

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- Motivation
- Dynamical Transport Equation
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Motivation

- Transport Simulation including Core and SOL Plasmas
 - Role of Separatrix
 - Closed magnetic surface ← Open magnetic field line
 - Difference of dominant transport process
- Tansient Behavior of Plasma Rotation
 - Radial Electric Field: Radial force balance, (Poisson equation)
 - Poloidal rotation: Equation of motion
 - Toroidal rotation: Equation of motion
 - Equation of motion rather than transport matrix
- Analysis including Atomic Processes

1D Transport code: TASK/TX

- Dynamic Transport Equation: Fukuyama et al. PPCF (1994)
 - Two fluid equation for electrons and ions
 - Flux surface averaged
 - Coupled with Maxwell equation
 - Neutral diffusion equation
 - Neoclassical transport
 - Turbulent transport
 - Current diffusive ballooning mode
 - Ambipolar diffusion through poloidal momentum transfer
 - Thermal diffusivity, Perpendicular viscosity
 - Maxwell's equation, Poisson's equation
 - Slowdown equation for beam component
 - Diffusion equation for neutral particles

Model Equation (1)

Fluid equations (electrons and ions)

$$\begin{split} \frac{\partial n_s}{\partial t} &= -\frac{1}{r} \frac{\partial}{\partial r} (r n_s u_{sr}) + S_s \\ \frac{\partial}{\partial t} (m_s n_s u_{sr}) &= -\frac{1}{r} \frac{\partial}{\partial r} (r m_s n_s u_{sr}^2) + \frac{1}{r} m_s n_s u_{s\theta}^2 + e_s n_s (E_r + u_{s\theta} B_{\phi} - u_{s\phi} B_{\theta}) - \frac{\partial}{\partial r} n_s T_s \\ \frac{\partial}{\partial t} (m_s n_s u_{s\theta}) &= -\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 m_s n_s u_{sr} u_{s\theta}) + e_s n_s (E_{\theta} - u_{sr} B_{\phi}) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^3 n_s m_s \mu_s \frac{\partial}{\partial r} \frac{u_{s\theta}}{r} \right) \\ &+ F_{s\theta}^{NC} + F_{s\theta}^{C} + F_{s\theta}^{NC} + F_{s\phi}^{NC} + F_{s\phi}^{NC}$$

Model Equation (2)

Diffusion equation for (fast and slow) neutral particles

$$\frac{\partial n_0}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} \left(-r D_0 \frac{\partial n_0}{\partial r} \right) + S_0$$

Maxwell's equation

$$\frac{1}{r}\frac{\partial}{\partial r}(rE_r) = \frac{1}{\epsilon_0}\sum_s e_s n_s$$

$$\frac{\partial B_{\theta}}{\partial t} = \frac{\partial E_{\phi}}{\partial r}, \qquad \frac{\partial B_{\phi}}{\partial t} = -\frac{1}{r}\frac{\partial}{\partial r}(rE_{\phi})$$

$$\frac{1}{c^2}\frac{\partial E_{\theta}}{\partial t} = -\frac{\partial}{\partial r}B_{\phi} - \mu_0\sum_s e_s n_s u_{s\theta}, \qquad \frac{1}{c^2}\frac{\partial E_{\phi}}{\partial t} = \frac{1}{r}\frac{\partial}{\partial r}(rB_{\theta}) - \mu_0\sum_s e_s n_s u_{s\phi}$$

Neoclassical Transport Model

Neoclassical transport

- Viscosity force arises when plasma rotates in the poloidal direction.
- Banana-Plateau regime

$$F_{S\theta}^{NC} = -\sqrt{\pi}q^2 n_S m_S \frac{v_{TS}}{qR} \frac{v_S^*}{1 + v_S^*} u_{S\theta}$$

$$v_S^* \equiv \frac{v_S qR}{\epsilon^{3/2} v_{TS}}$$

This poloidal viscosity force induces

- Neoclassical radial diffusion
- Neoclassical resistivity
- Bootstrap current
- Ware pinch

Turbulent Transport Model

Turbulent Diffusion

- Poloidal momentum exchange between electron and ion through the turbulent electric field
- Ambipolar flux (electron flux = ion flux)

$$F_{i\theta}^{W} = -F_{e\theta}^{W}$$

$$= -ZeB_{\phi}n_{\mathbf{i}}D_{\mathbf{i}}\left[-\frac{1}{n_{\mathbf{i}}}\frac{dn_{\mathbf{i}}}{dr} + \frac{Ze}{T_{\mathbf{i}}}E_{r} - \langle \frac{\omega}{m} \rangle \frac{ZeB_{\phi}}{T_{\mathbf{i}}} - \left(\frac{\mu_{\mathbf{i}}}{D_{\mathbf{i}}} - \frac{1}{2}\right)\frac{1}{T_{\mathbf{i}}}\frac{dT_{\mathbf{i}}}{dr}\right]$$

Perpendicular viscosity

- Non-ambipolar flux (electron flux \neq ion flux): $\mu_{S} = \text{constant} \times D$
- **Diffusion coefficient** (proportional to $|E|^2$)
 - Current-diffusive ballooning mode turbulence model

Model of Scrape-Off Layer Plasma

- Particle, momentum and heat losses along the field line
 - Decay time

$$\nu_{\rm L} = \begin{cases} 0 & (0 < r < a) \\ \frac{C_{\rm S}}{2\pi r R \{1 + \log[1 + 0.05/(r - a)]\}} & (a < r < b) \end{cases}$$

Electron source term

$$S_{\rm e} = n_0 \langle \sigma_{\rm ion} v \rangle n_{\rm e} - \nu_{\rm L} (n_{\rm e} - n_{\rm e, div})$$

- Recycling from divertor
 - \circ Recycling rate: $\gamma_0 = 0.8$
 - Neutral source

$$S_0 = \frac{\gamma_0}{Z_i} \nu_L (n_e - n_{e,div}) - \frac{1}{Z_i} n_0 \langle \sigma_{ion} v \rangle n_e + \frac{P_b}{E_b}$$

Gas puff from wall

Steady State Flux (1)

Electron flux

Inertia term in the equation of motion = 0

Radial flux

$$u_{er} = -\frac{1}{1 + \alpha} \frac{\bar{v}_{e} + v_{ei}}{n_{e} m_{e} \Omega_{e\phi}^{2}} \left[\frac{dP_{e}}{dr} + \frac{dP_{i}}{dr} \right] - \frac{\alpha}{1 + \alpha} \frac{E_{\phi}}{B_{\theta}} + \frac{1}{1 + \alpha} \frac{F_{\theta}^{W}}{n_{e} m_{e} \Omega_{e\phi}} + \frac{1}{1 + \alpha} \frac{A_{e\phi}}{n_{e} m_{e} \Omega_{e\phi}} \frac{B_{\phi} v_{eb}}{1 + \alpha} (u_{b\phi} - u_{i\phi})$$

where
$$\alpha \equiv \frac{\bar{\nu}_{\rm e} + \nu_{\rm ei}}{\nu_{\rm ei} + \nu_{\rm eb}} \frac{B_{\theta}^2}{B_{\phi}^2}$$
, $\nu_{\rm eb} = \frac{n_{\rm b} m_{\rm b}}{n_{\rm e} m_{\rm e}} \nu_{\rm be}$

- \circ Poloidal neoclassical viscosity $\bar{\nu}_e$
- \circ Factor α represents parallel neoclassical viscosity
- First three terms in RHS are neoclassical diffusion, Ware pinch and turbulent diffusion.

Steady State Flux (2)

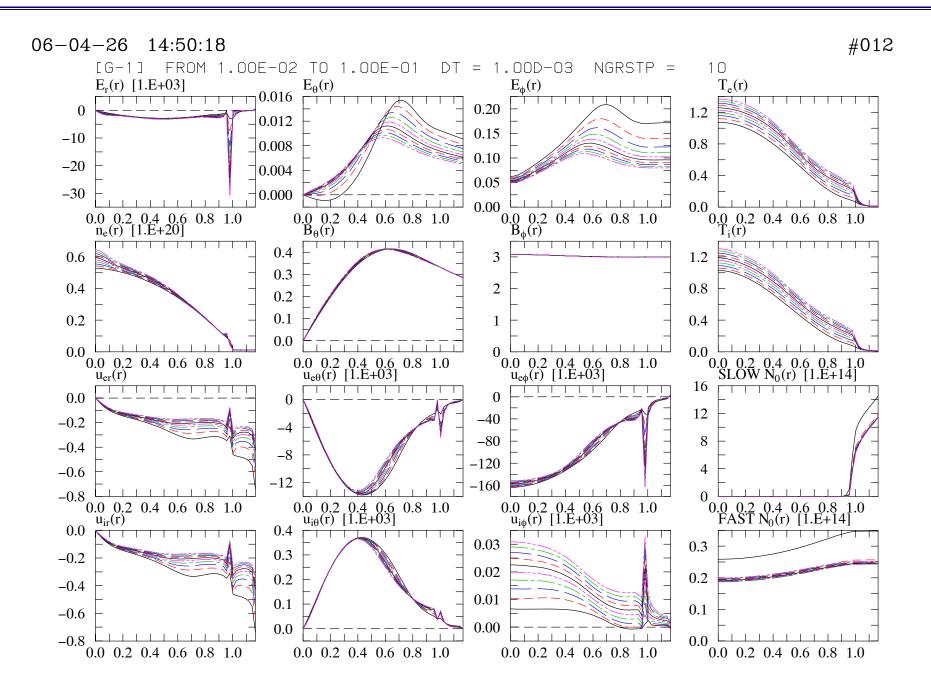
Toroidal current

$$u_{e\phi} = \frac{-1}{v_{ei} + v_{eb}} \left\{ \frac{1}{1 + \alpha} \frac{e}{m_e} E_{\phi} + \frac{1}{1 + \alpha} \frac{B_{\theta}}{B_{\phi}} \frac{\bar{v}_e + v_{ei}}{n_e m_e} \Omega_{e\phi} \left[\frac{dP_e}{dr} + \frac{dP_i}{dr} \right] \right.$$
$$\left. + \frac{1}{1 + \alpha} \frac{B_{\theta}}{B_{\phi}} \frac{F_{\theta}^{W}}{n_e m_e} + \frac{\bar{v}_e}{1 + \alpha} \frac{B_{\theta}}{B_{\phi}} u_{i\theta} \right.$$
$$\left. - \left(v_{eb} - \frac{\alpha}{1 + \alpha} v_{ei} \right) u_{b\phi} - \left(v_{ei} + \frac{\alpha}{1 + \alpha} v_{eb} \right) u_{i\phi} \right\}$$

- The first two terms in RHS are neoclassical resistivity and bootstrap current
- Similar expression for poloidal rotation

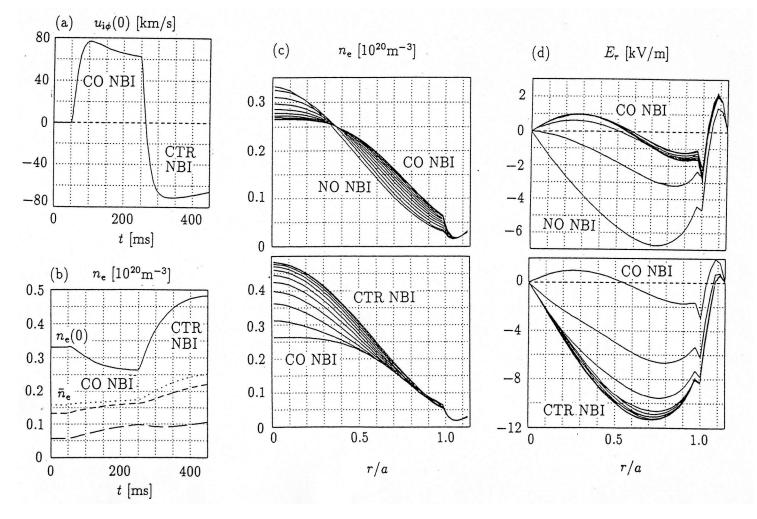
Model equations include dominant neoclassical transport.

Typical Profiles



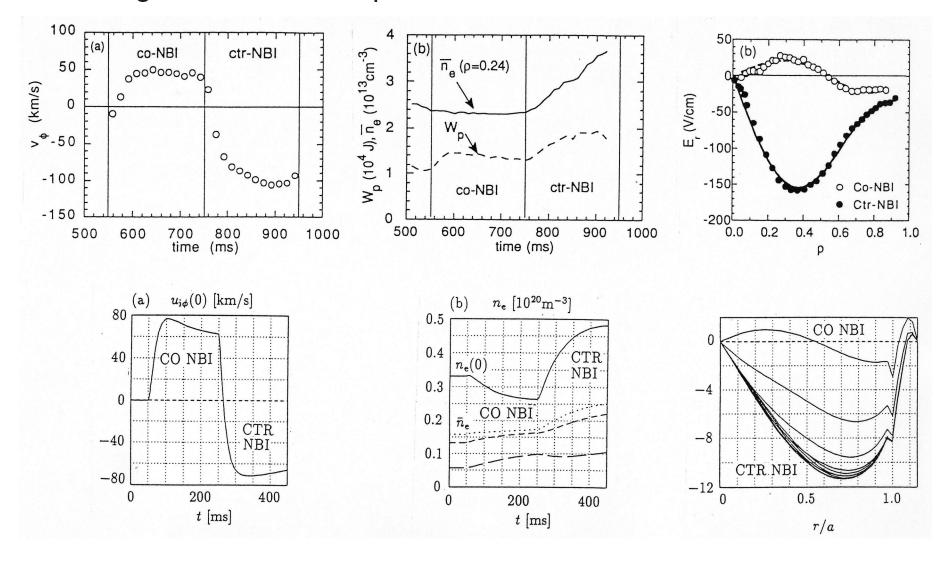
Simulation of plasma rotation and radial electric field

- JFT-2M parameter: NBI co-injection → counter-injection
- Toroidal rotation \implies Negative $E_r \implies$ Density peaking
- TASK/TX: Particle Diffusivity: $0.3 \,\mathrm{m}^2/\mathrm{s}$, lon viscosity: $10 \,\mathrm{m}^2/\mathrm{s}$



Comparison with JFT-2M Experiment

- JFT-2M Experiment: Ida et al.: Phys. Rev. Lett. 68 (1992) 182
- Good agreement with experimental observation



Density Peaking Due to Momentum Input

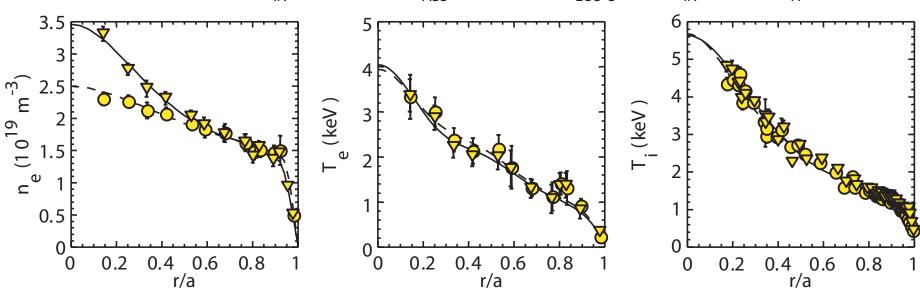
• Density peaking was observed in NBI counter injection on JT60-U. Ref. Takenaga et al. (ITPA, 2005)

oco-injection: $n_e(r/a=0.2)/< n_e>=1.41$

 P_{IN} =8.4MW, P_{ABS} =6.3MW, $P_{LOS S}$ FAST $/P_{IN}$ =0.26, H_{H} $^{98(y,2)}$ =0.86

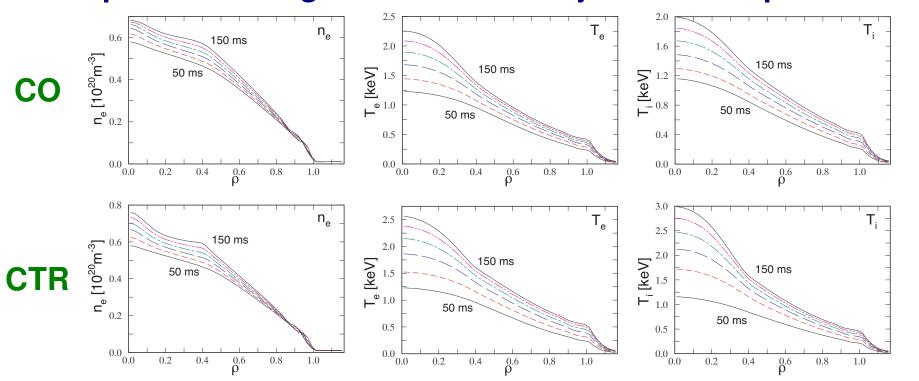
 ∇ ctr-injection: $n_e(r/a=0.2)/\langle n_e \rangle = 1.74$

 $P_{IN} = 10.2 \text{MW}, P_{ABS} = 6.5 \text{MW}, P_{LOS S}^{FAST} / P_{IN} = 0.38, H_{H}^{98(y,2)} = 0.81$



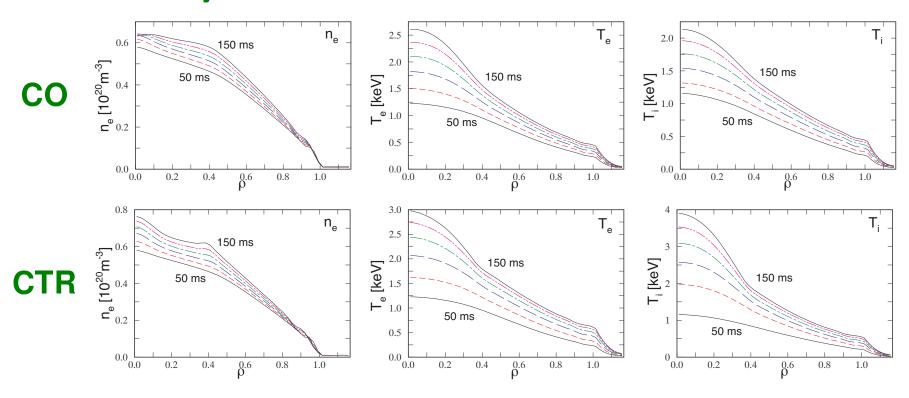
Density Peaking Simulation (1)

- Transport model: CDBM (particle diffusivity, thermal diffusivity)
- **NBI** 6.5 MW was injected 50 ms after the simulation started.
- Simulation results
 - \circ n(0) for co injection is 12% higher than that for counter injection
 - Temperature is higher for counter injection ↔ experiments



Density Peaking Simulation (2)

- NBI 10 MW heating
 - Co-injection: Density flattening
 - Counter-injection: ITB formation



Summary

- In order to describe the plasma rotation consistently, we have formulated a set of dynamical transport equations and implemented it a transport code TASK/TX.
- The numerical stability of TASK/TX was improved by using cubicspline FEM version.
- Density modification due to toroidal momentum input was successfully reproduced.
- Further improvement in modeling is required for realistic and robust simulation.