Modeling of Plasma Rotation and Density Modification

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Motivation

• Transport Simulation including Core and SOL Plasmas

• Role of Separatrix

— Closed magnetic surface ⇐⇒ Open magnetic field line
— Difference of dominant transport process

Tansient Behavior of Plasma Rotation

- Radial Electric Field: Radial force balance, (Poisson equation)
- **Poloidal rotation**: Equation of motion
- Toroidal rotation: Equation of motion
- Equation of motion rather than transport matrix
- Analysis including Atomic Processes

1D Transport code: TASK/TX

- **Dynamic Transport Equation**: *Fukuyama et al. PPCF (1994)*
 - Two fluid equation for electrons and ions
 - Flux surface averaged
 - Coupled with Maxwell equation
 - Neutral diffusion equation
 - Neoclassical transport
 - Turbulent transport
 - Current diffusive ballooning mode
 - Ambipolar diffusion through poloidal momentum transfer
 - Thermal diffusivity, Perpendicular viscosity
 - Maxwell's equation, Poisson's equation
 - Slowdown equation for beam component
 - Diffusion equation for neutral particles

• Fluid equations (electrons and ions)

$$\frac{\partial n_s}{\partial t} = -\frac{1}{r}\frac{\partial}{\partial r}\left(rn_s u_{sr}\right) + S_s$$

$$\frac{\partial}{\partial t}(m_s n_s u_{sr}) = -\frac{1}{r}\frac{\partial}{\partial r}(rm_s n_s u_{sr}^2) + \frac{1}{r}m_s n_s u_{s\theta}^2 + e_s n_s (E_r + u_{s\theta}B_{\phi} - u_{s\phi}B_{\theta}) - \frac{\partial}{\partial r}n_s T_s$$

$$\frac{\partial}{\partial t}(m_s n_s u_{s\theta}) = -\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 m_s n_s u_{sr} u_{s\theta}) + e_s n_s (E_\theta - u_{sr} B_\phi) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^3 n_s m_s \mu_s \frac{\partial}{\partial r} \frac{u_{s\theta}}{r} \right)$$

$$+F_{s\theta}^{\rm NC} + F_{s\theta}^{\rm C} + F_{s\theta}^{\rm W} + F_{s\theta}^{\rm X} + F_{s\theta}^{\rm L}$$

$$\frac{\partial}{\partial t} \left(m_s n_s u_{s\phi} \right) = -\frac{1}{r} \frac{\partial}{\partial r} (r m_s n_s u_{sr} u_{s\phi}) + e_s n_s (E_{\phi} + u_{sr} B_{\theta}) + \frac{1}{r} \frac{\partial}{\partial r} \left(r n_s m_s \mu_s \frac{\partial}{\partial r} u_{s\phi} \right)$$

$$+F_{s\phi}^{\mathrm{C}} + F_{s\phi}^{\mathrm{W}} + F_{s\phi}^{\mathrm{X}} + F_{s\phi}^{\mathrm{L}}$$

$$\frac{\partial}{\partial t}\frac{3}{2}n_{s}T_{s} = -\frac{1}{r}\frac{\partial}{\partial r}r\left(\frac{5}{2}u_{sr}n_{s}T_{s} - n_{s}\chi_{s}\frac{\partial}{\partial r}T_{e}\right) + e_{s}n_{s}(E_{\theta}u_{s\theta} + E_{\phi}u_{s\phi})$$
$$+P_{s}^{C} + P_{s}^{L} + P_{s}^{H}$$

• Diffusion equation for (fast and slow) neutral particles

$$\frac{\partial n_0}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} \left(-r D_0 \frac{\partial n_0}{\partial r} \right) + S_0$$

Maxwell's equation

$$\frac{1}{r}\frac{\partial}{\partial r}(rE_r) = \frac{1}{\epsilon_0}\sum_{s}e_s n_s$$
$$\frac{\partial B_{\theta}}{\partial t} = \frac{\partial E_{\phi}}{\partial r}, \qquad \frac{\partial B_{\phi}}{\partial t} = -\frac{1}{r}\frac{\partial}{\partial r}(rE_{\phi})$$
$$\frac{1}{c^2}\frac{\partial E_{\theta}}{\partial t} = -\frac{\partial}{\partial r}B_{\phi} - \mu_0\sum_{s}e_s n_s u_{s\theta}, \qquad \frac{1}{c^2}\frac{\partial E_{\phi}}{\partial t} = \frac{1}{r}\frac{\partial}{\partial r}(rB_{\theta}) - \mu_0\sum_{s}e_s n_s u_{s\phi}$$

Neoclassical Transport Model

Neoclassical transport

- Viscosity force arises when plasma rotates in the poloidal direction.
- Banana-Plateau regime

$$F_{s\theta}^{\rm NC} = -\sqrt{\pi}q^2 n_s m_s \frac{v_{\rm Ts}}{qR} \frac{v_s^*}{1 + v_s^*} u_{s\theta}$$
$$v_s^* \equiv \frac{v_s qR}{\epsilon^{3/2} v_{\rm Ts}}$$

- This poloidal viscosity force induces
 - Neoclassical radial diffusion
 - Neoclassical resistivity
 - Bootstrap current
 - Ware pinch

Turbulent Transport Model

Turbulent Diffusion

- Poloidal momentum exchange between electron and ion through the turbulent electric field
- Ambipolar flux (electron flux = ion flux)

 $F_{i\theta}^{W} = - F_{e\theta}^{W}$

$$= -ZeB_{\phi}n_{i}D_{i}\left[-\frac{1}{n_{i}}\frac{dn_{i}}{dr} + \frac{Ze}{T_{i}}E_{r} - \langle\frac{\omega}{m}\rangle\frac{ZeB_{\phi}}{T_{i}} - \left(\frac{\mu_{i}}{D_{i}} - \frac{1}{2}\right)\frac{1}{T_{i}}\frac{dT_{i}}{dr}\right]$$

Perpendicular viscosity

• Non-ambipolar flux (electron flux \neq ion flux): $\mu_s = \text{constant} \times D$

• **Diffusion coefficient** (proportional to $|E|^2$)

Current-diffusive ballooning mode turbulence model

• Particle, momentum and heat losses along the field line

• Decay time

$$v_{\rm L} = \begin{cases} 0 & (0 < r < a) \\ \frac{C_{\rm S}}{2\pi r R \{1 + \log[1 + 0.05/(r - a)]\}} & (a < r < b) \end{cases}$$

• Electron source term

$$S_{\rm e} = n_0 \langle \sigma_{\rm ion} v \rangle n_{\rm e} - \nu_{\rm L} (n_{\rm e} - n_{\rm e, div})$$

- Recycling from divertor
 - \circ Recycling rate: $\gamma_0 = 0.8$
 - Neutral source

$$S_0 = \frac{\gamma_0}{Z_i} v_L (n_e - n_{e,div}) - \frac{1}{Z_i} n_0 \langle \sigma_{ion} v \rangle n_e + \frac{P_b}{E_b}$$

Gas puff from wall

• Electron flux

 \circ Inertia term in the equation of motion = 0

Radial flux

$$\begin{split} u_{er} &= -\frac{1}{1+\alpha} \frac{\bar{v}_{e} + v_{ei}}{n_{e}m_{e}\Omega_{e\phi}^{2}} \left[\frac{\mathrm{d}P_{e}}{\mathrm{d}r} + \frac{\mathrm{d}P_{i}}{\mathrm{d}r} \right] - \frac{\alpha}{1+\alpha} \frac{E_{\phi}}{B_{\theta}} + \frac{1}{1+\alpha} \frac{F_{\theta}^{W}}{n_{e}m_{e}\Omega_{e\phi}} \\ &+ \frac{1}{1+\alpha} \frac{\bar{v}_{e}}{\Omega_{e\phi}} u_{i\theta} + \frac{\alpha}{1+\alpha} \frac{B_{\phi}}{B_{\theta}} \frac{v_{eb}}{\Omega_{e\phi}} (u_{b\phi} - u_{i\phi}) \\ &\text{where} \quad \alpha \equiv \frac{\bar{v}_{e} + v_{ei}}{v_{ei} + v_{eb}} \frac{B_{\theta}^{2}}{B_{\phi}^{2}}, \quad v_{eb} = \frac{n_{b}m_{b}}{n_{e}m_{e}} v_{be} \end{split}$$

 \circ Poloidal neoclassical viscosity $\bar{\nu}_e$

- \circ Factor α represents parallel neoclassical viscosity
- First three terms in RHS are neoclassical diffusion, Ware pinch and turbulent diffusion.

Toroidal current

$$u_{e\phi} = \frac{-1}{v_{ei} + v_{eb}} \left\{ \frac{1}{1 + \alpha} \frac{e}{m_e} E_{\phi} + \frac{1}{1 + \alpha} \frac{B_{\theta}}{B_{\phi}} \frac{\bar{v}_e + v_{ei}}{n_e m_e} \left[\frac{dP_e}{dr} + \frac{dP_i}{dr} \right] \right. \\ \left. + \frac{1}{1 + \alpha} \frac{B_{\theta}}{B_{\phi}} \frac{F_{\theta}^W}{n_e m_e} + \frac{\bar{v}_e}{1 + \alpha} \frac{B_{\theta}}{B_{\phi}} u_{i\theta} \right. \\ \left. - \left(v_{eb} - \frac{\alpha}{1 + \alpha} v_{ei} \right) u_{b\phi} - \left(v_{ei} + \frac{\alpha}{1 + \alpha} v_{eb} \right) u_{i\phi} \right\}$$

- The first two terms in RHS are neoclassical resistivity and bootstrap current
- Similar expression for poloidal rotation

Model equations include dominant neoclassical transport.

Typical Profiles



Simulation of plasma rotation and radial electric field

- JFT-2M parameter: NBI co-injection → counter-injection
- Toroidal rotation \implies Negative $E_r \implies$ Density peaking
- TASK/TX: Particle Diffusivity: $0.3 \text{ m}^2/\text{s}$, lon viscosity: $10 \text{ m}^2/\text{s}$



Comparison with JFT-2M Experiment

- JFT-2M Experiment: Ida et al.: Phys. Rev. Lett. 68 (1992) 182
- Good agreement with experimental observation



Density Peaking Due to Momentum Input

• Density peaking was observed in NBI counter injection on JT60-U. *Ref. Takenaga et al. (ITPA, 2005)*



Density Peaking Simulation (1)

- Transport model: CDBM (particle diffusivity, thermal diffusivity)
- **NBI** 6.5 MW was injected 50 ms after the simulation started.
- Simulation results
 - *n*(0) for co injection is 12% higher than that for counter injection
 - \circ Temperature is higher for counter injection \leftrightarrow experiments



Density Peaking Simulation (2)

• NBI 10 MW heating

• Co-injection: Density flattening

• Counter-injection: ITB formation



Summary

- In order to describe the plasma rotation consistently, we have formulated a set of dynamical transport equations and implemented it a transport code TASK/TX.
- The numerical stability of TASK/TX was improved by using cubicspline FEM version.
- Density modification due to toroidal momentum input was successfully reproduced.
- Further improvement in modeling is required for realistic and robust simulation.