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Dynamic Transport Simulation Including Plasma Rotation and Radial Electric Field

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- TASK/TX: Dynamic transport simulation code
- Physics included in the TASK/TX code
- Numerical results
- Summary

• Hierarchy of transport phenomena in toroidal plasmas:

- **TASK/TR**: Diffusive Transport Equations:
 - Gradient-flux relation: Stationary solution of eqs. of motion
 - Conventional way of transport simulations
- TASK/TX: Dynamic Transport Equations: ⇐ Main topic
 - Flux-surface averaged multi-fluid equations
 - Including inertia terms in equations of motion
 - Coupling with Maxwell's equations
 - Transient analysis including plasma rotations and E_r
- **TASK/FP**: Kinetic Transport Equations:
 - Bounce-averaged Fokker-Plank equations
 - Modification of momentum distribution functions
 - Integrating heating, current drive and kinetic stability analysis

Motivation of the TASK/TX Code

• Transport Simulation Including Core and SOL Plasmas

• Role of separatrix

- Closed magnetic surface \iff Open magnetic field line
- Difference of dominant transport processes
- Transient Behavior of Plasma Rotation
 - Radial electric field: Radial force balance (Gauss's law = Poisson's equation)
 - **Poloidal rotation**: Equation of motion
 - Toroidal rotation: Equation of motion
 - Equation of motion rather than transport matrix
- Analysis including Atomic Processes

1D Dynamic Transport Code: TASK/TX

- **Dynamic Transport Equations** (TASK/TX)
 - A set of flux-surface averaged equations
 - \circ Two fluid equations for electrons and ions
 - Continuity equations
 - Equations of motion (radial, poloidal and toroidal)
 - Energy transport equations
 - Neoclassical transport
 - Poloidal viscosity \implies Emerging all the neoclassical effects
 - Turbulent transport
 - Intrinsic ambipolar diffusion through poloidal momentum transf.
 - Thermal diffusivity and perpendicular viscosity
 - Maxwell's equations including Poisson's equation
 - Slowing-down equations for beam ion component
 - Diffusion equations for two-groups (fast and slow) neutrals

• Fluid equations for electrons and ions (s = e, i):

$$\frac{\partial n_s}{\partial t} = -\frac{1}{r}\frac{\partial}{\partial r}\left(rn_s u_{sr}\right) + S_s$$

$$\frac{\partial}{\partial t}(m_s n_s u_{sr}) = -\frac{1}{r}\frac{\partial}{\partial r}(rm_s n_s u_{sr}^2) + \frac{1}{r}m_s n_s u_{s\theta}^2 - \frac{\partial}{\partial r}n_s T_s + e_s n_s (E_r + u_{s\theta}B_{\phi} - u_{s\phi}B_{\theta})$$

$$\frac{\partial}{\partial t}(m_s n_s u_{s\theta}) = -\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 m_s n_s u_{sr} u_{s\theta}) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^3 n_s m_s \mu_s \frac{\partial}{\partial r} \frac{u_{s\theta}}{r}\right) + e_s n_s (E_\theta - u_{sr} B_\phi)$$

$$+ F_{s\theta}^{\rm NC} + F_{s\theta}^{\rm C} + F_{s\theta}^{\rm W} + F_{s\theta}^{\rm L} + F_{s\theta}^{\rm IN} + F_{s\theta}^{\rm CX}$$

$$\frac{\partial}{\partial t} \left(m_s n_s u_{s\phi} \right) = -\frac{1}{r} \frac{\partial}{\partial r} (r m_s n_s u_{sr} u_{s\phi}) + \frac{1}{r} \frac{\partial}{\partial r} \left(r n_s m_s \mu_s \frac{\partial}{\partial r} u_{s\phi} \right) + e_s n_s (E_\phi + u_{sr} B_\theta)$$

$$+ F^{\mathrm{C}}_{s\phi} + F^{\mathrm{W}}_{s\phi} + F^{\mathrm{L}}_{s\phi} + F^{\mathrm{IN}}_{s\phi} + F^{\mathrm{CX}}_{s\phi}$$

$$\frac{\partial}{\partial t}\frac{3}{2}n_sT_s = -\frac{1}{r}\frac{\partial}{\partial r}r\left(\frac{5}{2}u_{sr}n_sT_s - \frac{3}{2}n_s\chi_s\frac{\partial}{\partial r}T_e\right) + e_sn_s(E_\theta u_{s\theta} + E_\phi u_{s\phi})$$

 $+ P_s^{\rm C} + P_s^{\rm L} + P_s^{\rm R} + P_s^{\rm RF}$

Slowing-down equations for beam ion component

$$\frac{\partial n_{\rm b}}{\partial t} = S_{\rm b}^{B} - S_{\rm b}^{C}$$

$$\frac{\partial}{\partial t} \left(m_{\rm b} n_{\rm b} u_{\rm b\theta} \right) = e_{\rm b} n_{\rm b} E_{\theta} + F_{\rm b\theta}^{\rm C} + F_{\rm b\theta}^{\rm IN} + F_{\rm b\theta}^{\rm CX}$$

$$\frac{\partial}{\partial t} \left(m_{\rm b} n_{\rm b} u_{\rm b\phi} \right) = e_{\rm b} n_{\rm b} E_{\phi} + F_{\rm b\phi}^{\rm C} + F_{\rm b\phi}^{\rm IN} + F_{\rm b\phi}^{\rm CX} + F_{\rm b\phi}^{\rm B}$$

• Diffusion equations for two-group neutrals (fast and slow)

$$\frac{\partial n_0}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} \left(-r D_0 \frac{\partial n_0}{\partial r} \right) + S_0$$

Maxwell's equations

$$\frac{1}{r}\frac{\partial}{\partial r}(rE_r) = \frac{1}{\epsilon_0}\sum_{s}e_s n_s$$
$$\frac{\partial B_{\theta}}{\partial t} = \frac{\partial E_{\phi}}{\partial r}, \qquad \frac{\partial B_{\phi}}{\partial t} = -\frac{1}{r}\frac{\partial}{\partial r}(rE_{\theta})$$
$$\frac{1}{c^2}\frac{\partial E_{\theta}}{\partial t} = -\frac{\partial}{\partial r}B_{\phi} - \mu_0\sum_{s}e_s n_s u_{s\theta}, \qquad \frac{1}{c^2}\frac{\partial E_{\phi}}{\partial t} = \frac{1}{r}\frac{\partial}{\partial r}(rB_{\theta}) - \mu_0\sum_{s}e_s n_s u_{s\phi}$$

Neoclassical transport

Parallel viscous force due to a poloidal plasma rotation
 Valid for all three neoclassical regimes

$$F_{s\theta}^{\rm NC} \equiv -n_s m_s v_{\rm NCs} u_{s\theta} = -\frac{\langle B^2 \rangle \hat{\mu}_{11}^{si}}{n_s m_s B_{\theta}^2} n_s m_s u_{s\theta}$$

 $\hat{\mu}_{11}^{si}$: viscosity coefficient from the NCLASS module, W. A. Houlberg et al. PoP **4** (1997) 3230

• Due to the poloidal viscous force,

- Neoclassical diffusion and Ware pinch
- Resistivity and bootstrap current

Turbulent diffusion

- Poloidal momentum exchange between electrons and ions through turbulent fluctuating field
- Intrinsic ambipolar flux

(electron particle flux = ion particle flux)

$$F_{e\theta}^{W} = -F_{i\theta}^{W} = -\frac{e^2 B_{\phi}^2 D_e}{T_e} n_e \left(u_{e\theta} - \frac{B_{\theta}}{B_{\phi}} u_{e\phi} \right)$$
$$F_{e\phi}^{W} = -F_{i\phi}^{W} = \frac{e^2 B_{\phi}^2 D_e}{T_e} \frac{B_{\theta}}{B_{\phi}} n_e \left(u_{e\theta} - \frac{B_{\theta}}{B_{\phi}} u_{e\phi} \right)$$

Perpendicular viscosity

Non-ambipolar particle flux
 (electron particle flux ≠ ion particle flux)

Modelling of SOL Plasma

• Parallel losses in the SOL

• Particle, momentum and ion heat losses: convection

$$v_{\rm L} = \frac{k_{\rm L} C_{\rm s}}{2\pi q R} \quad (a < r < b)$$

• Electron heat loss: conduction

$$v_{\rm L} = k_{\rm L} \frac{\chi_{\parallel}}{(2\pi q R)^2} = k_{\rm L} \frac{\kappa_0 T_{\rm e}^{5/2}}{n_{\rm e} (2\pi q R)^2} \quad (a < r < b)$$

• Particle source

$$S_{\rm e} = n_0 \langle \sigma_{\rm ion} v \rangle n_{\rm e} - v_{\rm L} (n_{\rm e} - n_{\rm e, div})$$

Recycling from divertor

 \circ Recycling ratio: $\gamma_0 = 0.8$

Gas puff from wall

• Stationary Electron Flux

Setting inertia terms to zero in the model equations

Radial velocity

$$\begin{split} u_{\rm er} &= -\frac{1}{1+\alpha} \frac{\bar{v}_{\rm e} + v_{\rm eNC}}{n_{\rm e} m_{\rm e} \Omega_{\rm e\phi}^2} \frac{\mathrm{d}p}{\mathrm{d}r} - \frac{\alpha}{1+\alpha} \frac{E_{\phi}}{B_{\theta}} + \frac{1}{1+\alpha} \frac{1}{n_{\rm e} m_{\rm e} \Omega_{\rm e\phi}} \left(F_{\rm e\theta}^{\rm W} + \frac{B_{\phi}}{B_{\theta}} \alpha F_{\rm e\phi}^{\rm W}\right) \\ &+ \frac{\alpha}{1+\alpha} \frac{1}{\Omega_{\rm e\phi}} \frac{B_{\phi}}{B_{\theta}} \left[v_{\rm eb} u_{\rm b\phi} - (\bar{v}_{\rm e} - v_{\rm ei}) u_{\rm i\phi}\right] + \frac{1}{1+\alpha} \frac{\bar{v}_{\rm e} + v_{\rm eNC} - v_{\rm ei}}{\Omega_{\rm e\phi}} u_{\rm i\theta} \end{split}$$

where $\bar{\nu}_e \equiv \nu_{ei} + \nu_{eb} + \nu_L + \nu_{0e}$,

$$\alpha \equiv \frac{\bar{\nu}_{\rm e} + \nu_{\rm eNC}}{\bar{\nu}_{\rm e}} \frac{B_{\theta}^2}{B_{\phi}^2}, \quad \Omega_{\rm e\phi} \equiv \frac{eB_{\phi}}{m_{\rm e}}, \quad \text{and} \quad \nu_{\rm eb} \equiv \frac{n_{\rm b}m_{\rm b}}{n_{\rm e}m_{\rm e}} \nu_{\rm be}$$

 $^{\circ}$ Damping rate, $\bar{\nu}_{eNC},$ due to the neoclassical viscosity

- First, second and third terms in RHS are neoclassical diffusion,
 Ware pinch and turbulent diffusion, respectively.
- Fourth term denotes a neoclassical pinch due to a momentum input from beam ions.

Toroidal velocity

$$\begin{split} u_{e\phi} &= -\frac{1}{\bar{\nu}_{e}} \left[\frac{1}{1+\alpha} \frac{e}{m_{e}} E_{\phi} - \frac{1}{1+\alpha} \frac{B_{\theta}}{B_{\phi}} \frac{\bar{\nu}_{e} + \nu_{eNC}}{n_{e}m_{e}} \frac{dp}{dr} + \frac{1}{1+\alpha} \frac{1}{n_{e}m_{e}} \left(\frac{B_{\theta}}{B_{\phi}} F_{e\theta}^{W} - F_{e\phi}^{W} \right) - \frac{\nu_{eb}}{1+\alpha} u_{b\phi} \\ &+ \frac{1}{1+\alpha} \frac{B_{\theta}}{B_{\phi}} (\bar{\nu}_{e} + \nu_{eNC} - \nu_{ei}) u_{i\theta} - \frac{\nu_{ei} + \alpha \bar{\nu}_{e}}{1+\alpha} u_{i\phi} \right], \end{split}$$

- First, second and third terms in RHS are neoclassical resistivity, bootstrap current, and turbulent driven current.
- Poloidal velocity can be obtained in a similar way.

Model equations include major neoclassical effects!

Numerical Schemes Used in TASK/TX Code

• Finite element method (FEM)

- Linear interpolation function
- Streamline Upwind Petrov-Galerkin (SUPG) method
 - Stabilizing spurious oscillation due to first-derivative terms
- $\circ s = r^2$ coordinate rather than *r* coordinate
- Achieving higher mesh resolution near the separatrix

Time-advancing method

- Full-implicit method
 - Highly robust calculation
 - Time-consuming due to need of matrix equation solver
- Mass lumping method
- Picard method to solve nonlinear equations iteratively

Typical Ohmic Plasma Profiles at t = 50 ms

• JFT-2M like plasma composed of electron and hydrogen

 $R = 1.3 \text{ m}, a = 0.35 \text{ m}, b = 0.4 \text{ m}, B_{\phi b} = 1.3 \text{ T}, I_p = 0.2 \text{ MA}, S_{\text{puff}} = 5.0 \times 10^{18} \text{ m}^{-2} \text{s}^{-1}$ $\gamma = 0.8, Z_{\text{eff}} = 2.0$, Fixed turbulent coefficient profile



Neoclassical Transport without Turbulence

- Clarifying neoclassical transport
 - No turbulent diffusivity and viscosity
 - Fixed temperature profiles
- Density peaking with steep gradient near the separatrix
 - Density peaking due to Ware pinch
 - Inward flux in the SOL due to ionization of neutrals



Diffusion due to Turbulent Induced Force

- Confirming the validity of our particle diffusion model
 - No particle diffusion terms in the continuity equations

Turbulence-induced poloidal friction force \implies Change of radial velocity \implies Particle diffusion through convective term

- No neoclassical viscosity assumed in this case
- Obvious particle diffusion (Left) and no diffusion (Right)
- Particle diffusion desribed properly in our model.



Parameter Dependence on $D_{\rm e}$ and $I_{\rm p}$

- Profile modifications due to the change of particle diffusivity
 - \circ Density flattening with the increase of particle diffusion, D_e
 - Notch of E_r near the separatrix vanishes with the increase of D_e because of alleviation of $\frac{dp}{dr}$.



- Profile modifications due to the change of plasma current
 - \circ Increase in *n* and *T* near the axis and decrease in E_r with the increase of I_p



NBI of $P_{\rm NB} = 0.5$ MW at t = 100 ms

• The cases of before-, co- and ctr-NBIs



Comparison with JFT-2M experiment (1)

Ref. K. Ida et al., PRL 68 (1992) 182



Comparison with JFT-2M experiment (2)

Ref. K. Ida et al., PRL 68 (1992) 182



$D_{\rm e}$ dependence in the case of counter NBI

- Density profile is determined by the balance between turbulentdriven and neoclassical particle fluxes during NBI.
- Increase particle diffusivity from 0.01 to 0.03 for counter NBI

 $\circ D_e = 0.01$ (Black): **Density peaking**

 $\circ D_e = 0.02$ (Red): Almost the same profile before NBI

 $\circ D_e = 0.03$ (Blue): No peaking due to strong particle diffusion

 $\circ E_r$ is not significantly affected except near the separatrix.



Summary

- TASK/TX: We are developing the TASK/TX code in order to describe poloidal and toroidal rotations and radial electric field formation. The code simultaneously solves the two-fluid equations of motion; continuity equations and heat transport equations coupled with Maxwell's equations.
- **Transport modeling**: The validity of our approach of neoclassical transport and particle transport was confirmed.
- Ohmic: We have examined the *I*_p dependence of the density, temperature and radial electric field profiles.
- NBI: We have analyzed the modifications of density profiles during NBI in JFT-2M like plasma. Density flattening in the case of co-NBI and density peaking in the case of ctr-NBI are qualitatively reproduced.

Future Works

- Precise surface average including toroidal metrics
- Equation for heat flux to complete neoclassical transport
- Theory-based turbulent transport model
- Multi-spices ion transport
- Dynamic simulation of transport barrier formation
- Long time simulation with more numerical robustness